

ON VARIATIONAL ITERATION METHODS FOR THE FRACTIONAL MODEL OF BLOOD ETHANOL CONCENTRATION SYSTEM

G. MONZÓN

Instituto de Ciencias
Universidad Nacional de General Sarmiento
Buenos Aires
Argentina
e-mail: gmonzon@campus.ungs.edu.ar

Abstract

In this paper, we numerically approximate the solutions of the first-order fractional differential equation system that models blood ethanol concentration using a fractional variational iteration method (FVIM) developed for high-order differential equations. Comparative analysis with two different variational iteration methods, including one specifically tailored for this problem, shows that the FVIM achieves lower approximation errors.

1. Introduction

The use of fractional differential equations to model diverse phenomena has seen significant growth in recent years across numerous fields including finance, biology, mechanics and engineering [3, 4, 8-10, 17, 28, 36]. In particular, in chemical engineering, one such phenomenon is the

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concentration of ethanol in human blood. In [30], this process is analyzed using the following system of first-order fractional differential equations:

$$\begin{cases} D^\alpha v(t) = -\kappa^\alpha v(t), \\ D^\alpha u(t) = \kappa^\alpha v(t) - \eta^\alpha u(t), \\ v(0) = v_0, \quad u(0) = 0, \end{cases} \quad (1)$$

where $v(t)$ denotes the concentration of alcohol in the stomach at time t (mg/l), $u(t)$ represents the concentration of alcohol in the blood at time t (mg/l), κ and η are rate constants (min^{-1}), v_0 is the initial concentration of alcohol in the stomach (mg/l) and D^α denotes the Caputo fractional derivative of order $\alpha \in (0, 1)$. The study in [30] demonstrates that this fractional model offers a more accurate estimate of real data compared to models that use integer-order derivatives.

The exact solution of (1) can be obtained using the Laplace transform technique. Indeed, from [30], we know that

$$v(t) = v_0 E_\alpha(-\kappa^\alpha t^\alpha) \quad \text{and} \quad u(t) = v_0 \kappa^\alpha \sum_{i \geq 0} \sum_{j \geq 0} \frac{(-1)^{i+j} \kappa^{\alpha i} \eta^{\alpha j}}{\Gamma((i+j+1)\alpha + 1)} t^{(i+j+1)\alpha},$$

where E_α denotes the Mittag-Leffler function defined as

$$E_\alpha(z) = \sum_{i \geq 0} \frac{z^i}{\Gamma(i\alpha + 1)}.$$

However, having numerical methods that provide accurate and efficient approximations is highly useful and advantageous. Many numerical approximation techniques for fractional differential equations have been explored [2, 5, 7, 11-13, 15, 19-22, 26, 31, 33], among which the variational iteration method (VIM) stands out as one of the most popular and effective [16, 18, 25, 29, 32, 34, 35]. In particular, in [1], a VIM is developed exclusively for the problem (1). This method essentially relies on constructing each correction term using the Riemann-Liouville

operator and a Lagrange multiplier, which is determined through the stationary conditions. Moreover, the approximating sequences $(\phi_n)_n$ and $(\psi_n)_n$ for v and u , respectively, are given by

$$\phi_n(t) = \phi_{n-1}(t) - \frac{1}{\Gamma(\alpha)} \int_0^t (t-\tau)^{\alpha-1} (D^\alpha \phi_{n-1}(\tau) + \kappa^\alpha \phi_{n-1}(\tau)) d\tau, \quad (2a)$$

$$\begin{aligned} \psi_n(t) = \psi_{n-1}(t) - \frac{1}{\Gamma(\alpha)} \int_0^t (t-\tau)^{\alpha-1} (D^\alpha \psi_{n-1}(\tau) \\ - \kappa^\alpha \phi_{n-1}(\tau) + \eta^\alpha \psi_{n-1}(\tau)) d\tau, \end{aligned} \quad (2b)$$

with the initial terms

$$\phi_0(t) = v_0 \text{ and } \psi_0(t) = 0.$$

On the other hand, in [24], a VIM was proposed to obtain numerical approximations for arbitrary linear (and non-linear) systems of fractional differential equations. For the problem (1), the approximating sequences $(\varphi_n)_n$ and $(\chi_n)_n$ for v and u , respectively, are given by

$$\varphi_n(t) = \varphi_{n-1}(t) - \int_0^t (D^\alpha \varphi_{n-1}(\tau) + \kappa^\alpha \varphi_{n-1}(\tau)) d\tau, \quad (3a)$$

$$\chi_n(t) = \chi_{n-1}(t) - \int_0^t (D^\alpha \chi_{n-1}(\tau) - \kappa^\alpha \varphi_{n-1}(\tau) + \eta^\alpha \chi_{n-1}(\tau)) d\tau, \quad (3b)$$

with the initial terms

$$\varphi_0(t) = v_0 \text{ and } \chi_0(t) = 0.$$

In this work, we reformulate the partially coupled system (1) into a set of second-order homogeneous fractional differential equations and then, in order to obtain an alternative numerical method to (2) and (3), we use the fractional variational iteration method (FVIM) given in [27]. Through concrete numerical examples, we demonstrate that the fractional variational iteration method provides more accurate approximations.

Additionally, we derive alternative explicit expressions for the functions ϕ_n and ψ_n given in (2a) and (2b), which eliminate cross-term usage and unnecessary recurrences, relying solely on n (see Lemma 3.1 below). Clearly, this leads to a more straightforward and simplified implementation of the method. Nevertheless, the formulation we provide for the FVIM also offers this advantage and, as previously noted, yields improved results.

The organization of the paper is as follows: In Section 2, we introduce definitions, notations and basic properties related to the Riemann-Liouville integral operator, the Caputo fractional derivative and the Laplace transform, which are crucial for the comprehensive development of our work. Sections 3 and 4 are dedicated to presenting the approximations given by the variational iteration methods in [1] and [24], respectively, in a clear and concise manner. In Section 5, we use the FVIM developed in [27] to derive new approximating sequences for the real solutions of (1). Finally, in Section 6, numerical experiments are presented to compare the performance of the three methods considered.

2. Preliminaries

In this section, we briefly introduce definitions and well-known results that are useful for our purposes. For a more comprehensive and detailed discussion, we refer the reader to [6, 14, 23].

2.1. On the Riemann-Liouville integral operator and the Caputo fractional derivative

The Riemann-Liouville integral operator of w of order $\alpha \in (0, 1)$, $I^\alpha w$, is defined as

$$I^\alpha w(t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t - \tau)^{\alpha-1} w(\tau) d\tau \quad (t \geq 0), \quad (4)$$

where $\Gamma : (0, +\infty) \rightarrow \mathbb{R}$ denotes the Gamma function given by

$$\Gamma(x) = \int_0^{+\infty} \tau^{x-1} e^{-\tau} d\tau.$$

A well-known fact about this operator is that, for any $r \geq 0$,

$$I^\alpha t^r = \frac{\Gamma(r+1)}{\Gamma(r+\alpha+1)} t^{\alpha+r}. \tag{5}$$

On the other hand, the Caputo fractional derivative of w of order $\alpha \in (0, 1)$, $D^\alpha w$, can be defined by $I^{1-\alpha}$ as follows

$$D^\alpha w = I^{1-\alpha} w',$$

where w' denotes the ordinary derivative of w , i.e.,

$$D^\alpha w(t) = \frac{1}{\Gamma(1-\alpha)} \int_0^t \frac{w'(\tau)}{(t-\tau)^\alpha} d\tau \quad (t \geq 0). \tag{6}$$

Taking into account that Gamma function Γ satisfies

$$x\Gamma(x) = \Gamma(x+1),$$

from (5) it follows that, for any $r \geq 0$,

$$D^\alpha t^r = \frac{\Gamma(r+1)}{\Gamma(r+1-\alpha)} t^{r-\alpha}. \tag{7}$$

The operators I^α and D^α are also related in the following ways:

$$D^\alpha I^\alpha w = w, \tag{8}$$

and

$$I^\alpha D^\alpha w(t) = w(t) - w(0). \tag{9}$$

2.2. On the Laplace transform

The Laplace transform of v , Lv , is given by

$$Lv(s) = \int_0^{\infty} e^{-s\tau} v(\tau) d\tau \quad (s > 0). \quad (10)$$

Among the many properties of this linear operator, we emphasize the following:

$$Lt^r(s) = \frac{\Gamma(r+1)}{s^{r+1}}, \quad r \geq 0, \quad (11)$$

$$LI^\alpha v(s) = \frac{1}{s^\alpha} Lv(s), \quad (12)$$

and

$$LD^\alpha v(s) = s^\alpha Lv(s) - s^{\alpha-1}v(0). \quad (13)$$

Finally, as usual, we use L^{-1} to denote the inverse Laplace transform. In particular, from (11), L^{-1} is a linear operator which satisfies

$$L^{-1} \frac{1}{s^{r+1}}(t) = \frac{t^r}{\Gamma(r+1)}, \quad r \geq 0. \quad (14)$$

3. On the VIM developed in [1]

A variational iteration method for the problem (1) is exclusively developed in [1] where the correction terms are built from the Riemann-Liouville operator (4) and the Lagrange multipliers are obtained from the stationary conditions via variational theory. In fact, the proposed approximating sequences $(\phi_n)_n$ of v and $(\psi_n)_n$ of u (previously introduced here in (2a) and (2b)) can be actually written as

$$\phi_n = \phi_{n-1} - I^\alpha(D^\alpha\phi_{n-1} + \kappa^\alpha\phi_{n-1}), \quad \phi_0 = v_0,$$

and

$$\psi_n = \psi_{n-1} - I^\alpha(D^\alpha\psi_{n-1} - \kappa^\alpha\phi_{n-1} + \eta^\alpha\psi_{n-1}), \quad \psi_0 = 0.$$

Now, thanks to the linearity of the operator I^α and the property (9) it follows that

$$\phi_n = \phi_{n-1} - I^\alpha D^\alpha \phi_{n-1} - \kappa^\alpha I^\alpha \phi_{n-1} = v_0 - \kappa^\alpha I^\alpha \phi_{n-1}, \quad (15)$$

and then

$$\psi_n = \psi_{n-1} - I^\alpha D^\alpha \psi_{n-1} + \kappa^\alpha I^\alpha \phi_{n-1} - \eta^\alpha I^\alpha \psi_{n-1} = v_0 - \phi_n - \eta^\alpha I^\alpha \psi_{n-1}. \quad (16)$$

From these identities, we can derive alternative expressions for ϕ_n and ψ_n which are highly useful from a computational point of view since they depend solely on n , avoiding the cross-use of terms and unnecessary recurrences.

Lemma 3.1. *Let ϕ_n and ψ_n , with $n \geq 1$, be as in (15) and (16) respectively, then*

$$\phi_n(t) = v_0 \sum_{i=0}^n \frac{(-1)^i (\kappa t)^{\alpha i}}{\Gamma(i\alpha + 1)}, \quad (17)$$

and

$$\psi_n(t) = v_0 \kappa^\alpha \sum_{i=1}^n \frac{(-1)^{i+1} \left(\sum_{j=0}^{i-1} \kappa^{(i-1-j)\alpha} \eta^{j\alpha} \right) t^{i\alpha}}{\Gamma(i\alpha + 1)}. \quad (18)$$

Proof. We argue by induction on n . For $n = 1$, from (15) and (5) we have

$$\begin{aligned} \phi_1(t) &= v_0 - \kappa^\alpha I^\alpha v_0 = v_0 (1 - \kappa^\alpha I^\alpha 1) = v_0 \left(1 - \kappa^\alpha \frac{t^\alpha}{\Gamma(\alpha + 1)} \right) \\ &= v_0 \sum_{i=0}^1 \frac{(-1)^i (\kappa^\alpha t^\alpha)^i}{\Gamma(i\alpha + 1)}; \end{aligned}$$

on the other hand, from (16) we have

$$\psi_1(t) = v_0 - \phi_1(t) - \eta^\alpha I^\alpha 0 = v_0 - \phi_1(t) = v_0 \kappa^\alpha \frac{t^\alpha}{\Gamma(\alpha + 1)},$$

which proves the claim.

Now suppose that (17) holds for $n \geq 1$. From (15) and (5), it follows that

$$\begin{aligned} \phi_{n+1}(t) &= v_0 - \kappa^\alpha I^\alpha \phi_n(t) \\ &= v_0 - \kappa^\alpha I^\alpha \left(v_0 \sum_{i=0}^n \frac{(-1)^i \kappa^{\alpha i}}{\Gamma(i\alpha + 1)} t^{\alpha i} \right) \\ &= v_0 - \kappa^\alpha v_0 \sum_{i=0}^n \frac{(-1)^i \kappa^{\alpha i}}{\Gamma(i\alpha + 1)} I^\alpha t^{\alpha i} \\ &= v_0 + v_0 \sum_{i=0}^n \frac{(-1)^{i+1} \kappa^{\alpha(i+1)}}{\Gamma(i\alpha + 1)} \frac{\Gamma(i\alpha + 1) t^{\alpha(i+1)}}{\Gamma((i+1)\alpha + 1)} \\ &= v_0 \left(1 + \sum_{i=0}^n \frac{(-1)^{i+1} \kappa^{\alpha(i+1)} t^{\alpha(i+1)}}{\Gamma((i+1)\alpha + 1)} \right) = v_0 \sum_{i=0}^{n+1} \frac{(-1)^i (\kappa t)^{\alpha i}}{\Gamma(i\alpha + 1)}, \end{aligned}$$

which completes the argument regarding ϕ_n .

On the other hand, suppose that (18) holds for $n \geq 1$. From (16) and (5), it follows that

$$\begin{aligned} \psi_{n+1}(t) &= v_0 - \phi_{n+1}(t) - \eta^\alpha I^\alpha \psi_n(t) \\ &= v_0 - v_0 \sum_{i=0}^{n+1} \frac{(-1)^i (\kappa t)^{\alpha i}}{\Gamma(i\alpha + 1)} - \eta^\alpha I^\alpha \left(v_0 \kappa^\alpha \sum_{i=1}^n \frac{(-1)^{i+1} \left(\sum_{j=0}^{i-1} \kappa^{(i-1-j)\alpha} \eta^{j\alpha} \right) t^{\alpha i}}{\Gamma(i\alpha + 1)} \right) \\ &= v_0 \sum_{i=1}^{n+1} \frac{(-1)^{i+1} (\kappa t)^{\alpha i}}{\Gamma(i\alpha + 1)} - v_0 \eta^\alpha \kappa^\alpha \sum_{i=1}^n \frac{(-1)^{i+1} \left(\sum_{j=0}^{i-1} \kappa^{(i-1-j)\alpha} \eta^{j\alpha} \right)}{\Gamma(i\alpha + 1)} I^\alpha t^{\alpha i} \end{aligned}$$

$$\begin{aligned}
 &= v_0 \sum_{i=1}^{n+1} \frac{(-1)^{i+1} (\kappa t)^{\alpha i}}{\Gamma(i\alpha + 1)} - v_0 \eta^\alpha \kappa^\alpha \sum_{i=1}^n \frac{(-1)^{i+1} \left(\sum_{j=0}^{i-1} \kappa^{(i-1-j)\alpha} \eta^{j\alpha} \right) t^{\alpha i}}{\Gamma(i\alpha + 1)} \\
 &\quad \times \frac{\Gamma(i\alpha + 1) t^{\alpha(i+1)}}{\Gamma((i+1)\alpha + 1)} \\
 &= v_0 \sum_{i=1}^{n+1} \frac{(-1)^{i+1} (\kappa t)^{\alpha i}}{\Gamma(i\alpha + 1)} + v_0 \eta^\alpha \kappa^\alpha \sum_{i=2}^{n+1} \frac{(-1)^{i+1} \left(\sum_{j=0}^{i-2} \kappa^{(i-2-j)\alpha} \eta^{j\alpha} \right) t^{\alpha i}}{\Gamma(i\alpha + 1)} \\
 &= v_0 \kappa^\alpha \left[\begin{aligned} &\frac{t^\alpha}{\Gamma(\alpha + 1)} + \sum_{i=2}^{n+1} \frac{(-1)^{i+1} \kappa^{\alpha(i-1)} t^{\alpha i}}{\Gamma(i\alpha + 1)} \\ &+ \sum_{i=2}^{n+1} \frac{(-1)^{i+1} \left(\sum_{j=0}^{i-2} \kappa^{(i-2-j)\alpha} \eta^{(j+1)\alpha} \right) t^{\alpha i}}{\Gamma(i\alpha + 1)} \end{aligned} \right] \\
 &= v_0 \kappa^\alpha \left[\frac{t^\alpha}{\Gamma(\alpha + 1)} + \sum_{i=2}^{n+1} \frac{(-1)^{i+1} t^{\alpha i}}{\Gamma(i\alpha + 1)} \left(\kappa^{\alpha(i-1)} + \sum_{j=0}^{i-2} \kappa^{(i-2-j)\alpha} \eta^{(j+1)\alpha} \right) \right] \\
 &= v_0 \kappa^\alpha \left[\frac{t^\alpha}{\Gamma(\alpha + 1)} + \sum_{i=2}^{n+1} \frac{(-1)^{i+1} t^{\alpha i}}{\Gamma(i\alpha + 1)} \left(\sum_{j=0}^{i-1} \kappa^{(i-1-j)\alpha} \eta^{j\alpha} \right) \right] \\
 &= v_0 \kappa^\alpha \sum_{i=1}^{n-1} \frac{(-1)^{i+1} t^{\alpha i}}{\Gamma(i\alpha + 1)} \left(\sum_{j=0}^{i-1} \kappa^{(i-1-j)\alpha} \eta^{j\alpha} \right).
 \end{aligned}$$

This proves the lemma. \square

Remark 3.1. The explicit expression for ψ_3 given in [1] is

$$\psi_3(t) = v_0 \kappa^\alpha \left[\begin{aligned} &\frac{t^\alpha}{\Gamma(\alpha + 1)} - \frac{4^{-\alpha} (\kappa^\alpha + \eta^\alpha) \sqrt{\pi} t^{2\alpha}}{\Gamma(0.5 + \alpha) \Gamma(1 + \alpha)} \\ &+ \frac{4^{-\alpha} (\kappa^{2\alpha} + (\eta\kappa)^\alpha + \eta^{2\alpha}) \sqrt{\pi} \Gamma(1 + 2\alpha) t^{3\alpha}}{\Gamma(0.5 + \alpha) \Gamma(1 + \alpha) \Gamma(1 + 3\alpha)} \end{aligned} \right],$$

whereas the expression derived from (18) is

$$\psi_3(t) = v_0 \kappa^\alpha \left[\frac{t^\alpha}{\Gamma(\alpha + 1)} - \frac{(\kappa^\alpha + \eta^\alpha) t^{2\alpha}}{\Gamma(2\alpha + 1)} + \frac{(\kappa^{2\alpha} + (\eta\kappa)^\alpha + \eta^{2\alpha}) t^{3\alpha}}{\Gamma(3\alpha + 1)} \right].$$

The differences between both expressions can be explained using the Legendre's duplication formula:

$$\Gamma(\beta)\Gamma(0.5 + \beta) = 2^{1-2\beta}\sqrt{\pi}\Gamma(2\beta),$$

in the particular case $\beta = 0.5 + \alpha$. A similar situation arises with the expressions for ϑ_2 , ϕ_2 and ϕ_3 provided in [1] compared to those derived here.

4. On the VIM developed in [24]

Unlike the VIM proposed in [1], it is neither simple nor straightforward to find expressions for the approximating sequences φ_n and χ_n provided in [24] (given here in (3a)-(3b)) that depend solely on n . We present only the first five terms of each sequence, which will suffice for our purposes:

$$\varphi_1(t) = v_0(1 - \kappa^\alpha t),$$

$$\varphi_2(t) = v_0 \left(1 - 2\kappa^\alpha t + \frac{\kappa^{2\alpha}}{2} t^2 + \frac{\kappa^\alpha}{\Gamma(3-\alpha)} t^{2-\alpha} \right),$$

$$\varphi_3(t) = v_0 \left(1 - 3\kappa^\alpha t + \frac{3\kappa^{2\alpha}}{2} t^2 - \frac{\kappa^{3\alpha}}{6} t^3 + \frac{3\kappa^\alpha}{\Gamma(3-\alpha)} t^{2-\alpha} \right. \\ \left. - \frac{2\kappa^{2\alpha}}{\Gamma(4-\alpha)} t^{3-\alpha} - \frac{\kappa^\alpha}{\Gamma(4-2\alpha)} t^{3-2\alpha} \right),$$

$$\varphi_4(t) = v_0 \left(1 - 4\kappa^\alpha t + 3\kappa^{2\alpha} t^2 - \frac{2\kappa^{3\alpha}}{3} t^3 + \frac{\kappa^{4\alpha}}{24} t^4 + \frac{6\kappa^\alpha}{\Gamma(3-\alpha)} t^{2-\alpha} - \frac{8\kappa^{2\alpha}}{\Gamma(4-\alpha)} t^{3-\alpha} \right. \\ \left. - \frac{4\kappa^\alpha}{\Gamma(4-2\alpha)} t^{3-2\alpha} + \frac{3\kappa^{3\alpha}}{\Gamma(5-\alpha)} t^{4-\alpha} + \frac{3\kappa^{2\alpha}}{\Gamma(5-2\alpha)} t^{4-2\alpha} + \frac{\kappa^\alpha}{\Gamma(5-3\alpha)} t^{4-3\alpha} \right),$$

$$\begin{aligned}
 \varphi_5(t) = v_0 & \left(1 - 5\kappa^\alpha t + 5\kappa^{2\alpha} t^2 - \frac{5\kappa^{3\alpha}}{3} t^3 + \frac{5\kappa^{4\alpha}}{24} t^4 - \frac{\kappa^{5\alpha}}{120} t^5 \right. \\
 & + \frac{10\kappa^\alpha}{\Gamma(3-\alpha)} t^{2-\alpha} - \frac{20\kappa^{2\alpha}}{\Gamma(4-\alpha)} t^{3-\alpha} - \frac{10\kappa^\alpha}{\Gamma(4-2\alpha)} t^{3-2\alpha} + \frac{15\kappa^{3\alpha}}{\Gamma(5-\alpha)} t^{4-\alpha} \\
 & + \frac{15\kappa^{2\alpha}}{\Gamma(5-2\alpha)} t^{4-2\alpha} + \frac{5\kappa^\alpha}{\Gamma(5-3\alpha)} t^{4-3\alpha} - \frac{4\kappa^{4\alpha}}{\Gamma(6-\alpha)} t^{5-\alpha} \\
 & \left. - \frac{6\kappa^{3\alpha}}{\Gamma(6-2\alpha)} t^{5-2\alpha} - \frac{4\kappa^{2\alpha}}{\Gamma(6-3\alpha)} t^{5-3\alpha} - \frac{\kappa^\alpha}{\Gamma(6-4\alpha)} t^{5-4\alpha} \right),
 \end{aligned}$$

and, calling $\sigma_i = \sum_{j=0}^i \kappa^{(i-j)\alpha} \eta^{j\alpha}$ for $i = 1, \dots, 4$, we have

$$\chi_1(t) = v_0 \kappa^\alpha t,$$

$$\chi_2(t) = v_0 \kappa^\alpha \left(2t - \frac{\sigma_1}{2} t^2 - \frac{1}{\Gamma(3-\alpha)} t^{2-\alpha} \right),$$

$$\begin{aligned}
 \chi_3(t) = v_0 \kappa^\alpha & \left(3t - \frac{3\sigma_1}{2} t^2 + \frac{\sigma_2}{6} t^3 - \frac{3}{\Gamma(3-\alpha)} t^{2-\alpha} \right. \\
 & \left. + \frac{2\sigma_1}{\Gamma(4-\alpha)} t^{3-\alpha} + \frac{1}{\Gamma(4-2\alpha)} t^{3-2\alpha} \right),
 \end{aligned}$$

$$\begin{aligned}
 \chi_4(t) = v_0 \kappa^\alpha & \left(4t - 3\sigma_1 t^2 + \frac{2\sigma_2}{3} t^3 - \frac{\sigma_3}{24} t^4 - \frac{6}{\Gamma(3-\alpha)} t^{2-\alpha} + \frac{8\sigma_1}{\Gamma(4-\alpha)} t^{3-\alpha} \right. \\
 & \left. + \frac{4}{\Gamma(4-2\alpha)} t^{3-2\alpha} - \frac{3\sigma_2}{\Gamma(5-\alpha)} t^{4-\alpha} - \frac{3\sigma_1}{\Gamma(5-2\alpha)} t^{4-2\alpha} - \frac{1}{\Gamma(5-3\alpha)} t^{4-3\alpha} \right),
 \end{aligned}$$

$$\begin{aligned}
\chi_5(t) = & v_0 \kappa^\alpha \left(5t - 5\sigma_1 t^2 + \frac{5\sigma_2}{3} t^3 - \frac{5\sigma_3}{24} t^4 + \frac{\sigma_4}{120} t^5 - \frac{10}{\Gamma(3-\alpha)} t^{2-\alpha} \right. \\
& + \frac{20\sigma_1}{\Gamma(4-\alpha)} t^{3-\alpha} + \frac{10}{\Gamma(4-2\alpha)} t^{3-2\alpha} - \frac{15\sigma_2}{\Gamma(5-\alpha)} t^{4-\alpha} - \frac{15\sigma_1}{\Gamma(5-2\alpha)} t^{4-2\alpha} \\
& - \frac{5}{\Gamma(5-3\alpha)} t^{4-3\alpha} + \frac{4\sigma_3}{\Gamma(6-\alpha)} t^{5-\alpha} + \frac{6\sigma_2}{\Gamma(6-2\alpha)} t^{5-2\alpha} \\
& \left. + \frac{4\sigma_1}{\Gamma(6-3\alpha)} t^{5-3\alpha} + \frac{1}{\Gamma(6-4\alpha)} t^{5-4\alpha} \right).
\end{aligned}$$

5. Application of the FVIM [27]

From the partially coupled system (1), we can derive the following second-order fractional differential equation for the function u :

$$\begin{cases} D^{2\alpha} u + (\kappa^\alpha + \eta^\alpha) D^\alpha u + (\kappa\eta)^\alpha u = 0, \\ u(0) = 0, \quad D^\alpha u(0) = v_0 \kappa^\alpha, \end{cases}$$

where $D^{2\alpha}$ denotes the iterated derivative $D^\alpha \circ D^\alpha$. Clearly, this equation can be rewritten as

$$\begin{cases} D^{2\alpha} u = p D^\alpha u + q u, \\ u(0) = 0, \quad D^\alpha u(0) = v_0 \kappa^\alpha, \end{cases} \quad (19)$$

where

$$p = -(\kappa^\alpha + \eta^\alpha) \text{ and } q = -(\kappa\eta)^\alpha.$$

Applying the fractional variational iteration method proposed in [27], we obtain the approximating sequence $(u_n)_n$ given by

$$u_n(t) = \frac{v_0 \kappa^\alpha}{\Gamma(\alpha+1)} t^\alpha - L^{-1} \left[\frac{((\kappa^\alpha + \eta^\alpha) s^\alpha + (\kappa\eta)^\alpha) L u_{n-1}}{s^{2\alpha}} \right] \quad (n \geq 1), \quad (20)$$

with the initial data

$$u_0(t) = \frac{v_0 \kappa^\alpha}{\Gamma(\alpha + 1)} t^\alpha.$$

On the other hand, in order to obtain numerical approximations of the function v , it may be natural to consider the following single-term differential equation:

$$\begin{cases} D^\alpha v(t) = -\kappa^\alpha v(t), \\ u(0) = v_0. \end{cases}$$

However, we show that the FVIM proposed in [27] can also be applied in this case (and more generally to any first-order single-term fractional differential equation) with highly favorable results. Since a high-order fractional differential equation is necessary to apply such a method, we observe that the following second-order fractional differential equation can be derived from the system (1):

$$\begin{cases} D^{2\alpha} v(t) = \kappa^{2\alpha} v(t), \\ u(0) = v_0, \quad D^\alpha v(0) = -\kappa^\alpha v_0. \end{cases} \quad (21)$$

Thus, the approximating sequence provided by the FVIM in [27] is given by

$$v_n(t) = v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + \kappa^{2\alpha} L^{-1} \left[\frac{1}{s^{2\alpha}} L v_{n-1} \right] \quad (n \geq 1), \quad (22)$$

with the initial data

$$v_0(t) = v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right).$$

Lemma 5.1. *Let u_n and v_n , with $n \geq 1$, be as in (20) and (22) respectively, then*

$$u_n(t) = v_0 \kappa^\alpha \sum_{j=0}^n \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa \eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha + 1)} \quad (23)$$

and

$$v_n(t) = v_0 \sum_{i=0}^{2n+1} \frac{(-\kappa^\alpha t^\alpha)^i}{\Gamma(i\alpha + 1)}. \quad (24)$$

Proof. We argue by induction on n . For $n = 1$, from (20), (11) and (14) we have

$$\begin{aligned} u_1(t) &= \frac{v_0 \kappa^\alpha}{\Gamma(\alpha + 1)} t^\alpha - L^{-1} \left[\frac{((\kappa^\alpha + \eta^\alpha) s^\alpha + (\kappa\eta)^\alpha) L u_0(s)}{s^{2\alpha}} \right] \\ &= \frac{v_0 \kappa^\alpha}{\Gamma(\alpha + 1)} t^\alpha - L^{-1} \left[\frac{((\kappa^\alpha + \eta^\alpha) s^\alpha + (\kappa\eta)^\alpha) L \left[\frac{v_0 \kappa^\alpha}{\Gamma(\alpha + 1)} t^\alpha \right]}{s^{2\alpha}} \right] \\ &= \frac{v_0 \kappa^\alpha}{\Gamma(\alpha + 1)} t^\alpha - L^{-1} \left[\frac{((\kappa^\alpha + \eta^\alpha) s^\alpha + (\kappa\eta)^\alpha) \frac{v_0 \kappa^\alpha}{s^{\alpha+1}}}{s^{2\alpha}} \right] \\ &= v_0 \kappa^\alpha \left(\frac{t^\alpha}{\Gamma(\alpha + 1)} - (\kappa^\alpha + \eta^\alpha) L^{-1} \left[\frac{1}{s^{2\alpha+1}} \right] - (\kappa\eta)^\alpha L^{-1} \left[\frac{1}{s^{3\alpha+1}} \right] \right) \\ &= v_0 \kappa^\alpha \left(\frac{t^\alpha}{\Gamma(\alpha + 1)} - (\kappa^\alpha + \eta^\alpha) \frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} - (\kappa\eta)^\alpha \frac{t^{3\alpha}}{\Gamma(3\alpha + 1)} \right) \\ &= v_0 \kappa^\alpha \sum_{j=0}^1 \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha + 1)}; \end{aligned}$$

and, from (22), (11) and (14) it follows that

$$\begin{aligned} v_1(t) &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + \kappa^{2\alpha} L^{-1} \left[\frac{1}{s^{2\alpha}} L v_0(s) \right] \\ &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + \kappa^{2\alpha} L^{-1} \left[\frac{1}{s^{2\alpha}} L \left[v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) \right] \right] \end{aligned}$$

$$\begin{aligned}
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + \kappa^{2\alpha} L^{-1} \left[\frac{v_0}{s^{2\alpha}} \left(\frac{1}{s} - \frac{\kappa^\alpha}{s^{\alpha+1}} \right) \right] \\
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} + \kappa^{2\alpha} \left(L^{-1} \left[\frac{1}{s^{2\alpha+1}} \right] - \kappa^\alpha L^{-1} \left[\frac{1}{s^{3\alpha+1}} \right] \right) \right) \\
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} + \kappa^{2\alpha} \left(\frac{t^{2\alpha}}{\Gamma(2\alpha + 1)} - \frac{\kappa^\alpha t^{3\alpha}}{\Gamma(3\alpha + 1)} \right) \right) \\
 &= v_0 \sum_{i=0}^3 \frac{(-\kappa^\alpha t^\alpha)^i}{\Gamma(i\alpha + 1)},
 \end{aligned}$$

which proves the claim.

Now, suppose that (23) holds for $n \geq 1$. Then, thanks to (11),

$$\begin{aligned}
 Lu_n(s) &= L \left[v_0 \kappa^\alpha \sum_{j=0}^n \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha + 1)} \right] \\
 &= v_0 \kappa^\alpha \sum_{j=0}^n \sum_{i=0}^j \binom{j}{i} (-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} L \left[\frac{t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha + 1)} \right] \\
 &= v_0 \kappa^\alpha \sum_{j=0}^n \sum_{i=0}^j \binom{j}{i} (-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} \frac{1}{s^{(i+j+1)\alpha+1}}.
 \end{aligned}$$

From (20), the newly derived identity and (14), it follows that

$$\begin{aligned}
 u_{n+1}(t) &= \frac{v_0 \kappa^\alpha}{\Gamma(\alpha + 1)} t^\alpha - L^{-1} \left[\frac{((\kappa^\alpha + \eta^\alpha) s^\alpha + (\kappa\eta)^\alpha) Lu_n(s)}{s^{2\alpha}} \right] \\
 &= \frac{v_0 \kappa^\alpha}{\Gamma(\alpha + 1)} t^\alpha - v_0 \kappa^\alpha L^{-1} \left[\sum_{j=0}^n \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i+1} (\kappa\eta)^{i\alpha}}{s^{(i+j+2)\alpha+1}} \right. \\
 &\quad \left. + \sum_{j=0}^n \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{(i+1)\alpha}}{s^{(i+j+3)\alpha+1}} \right]
 \end{aligned}$$

$$\begin{aligned}
&= v_0 \kappa^\alpha \left(\frac{t^\alpha}{\Gamma(\alpha+1)} - \sum_{j=0}^n \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i+1} (\kappa\eta)^{i\alpha} t^{(i+j+2)\alpha}}{\Gamma((i+j+2)\alpha+1)} \right. \\
&\quad \left. - \sum_{j=0}^n \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{(i+1)\alpha} t^{(i+j+3)\alpha}}{\Gamma((i+j+3)\alpha+1)} \right) \\
&= v_0 \kappa^\alpha \left(\frac{t^\alpha}{\Gamma(\alpha+1)} + \sum_{j=1}^{n+1} \sum_{i=0}^{j-1} \binom{j-1}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha+1)} \right. \\
&\quad \left. + \sum_{j=1}^{n+1} \sum_{i=1}^j \binom{j-1}{i-1} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha+1)} \right) \\
&= v_0 \kappa^\alpha \left(\frac{t^\alpha}{\Gamma(\alpha+1)} + \sum_{j=1}^{n+1} \left(\sum_{i=1}^{j-1} \left(\binom{j-1}{i} + \binom{j-1}{i-1} \right) \right. \right. \\
&\quad \left. \left. \times \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha+1)} \right. \right. \\
&\quad \left. \left. + \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^j t^{(j+1)\alpha}}{\Gamma((j+1)\alpha+1)} + \frac{(-1)^j (\kappa\eta)^{j\alpha} t^{(2j+1)\alpha}}{\Gamma((2j+1)\alpha+1)} \right) \right) \\
&= v_0 \kappa^\alpha \left(\frac{t^\alpha}{\Gamma(\alpha+1)} + \sum_{j=1}^{n+1} \left(\sum_{i=1}^{j-1} \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha+1)} \right. \right. \\
&\quad \left. \left. + \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^j t^{(j+1)\alpha}}{\Gamma((j+1)\alpha+1)} + \frac{(-1)^j (\kappa\eta)^{j\alpha} t^{(2j+1)\alpha}}{\Gamma((2j+1)\alpha+1)} \right) \right) \\
&= v_0 \kappa^\alpha \left(\frac{t^\alpha}{\Gamma(\alpha+1)} + \sum_{j=1}^{n+1} \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha+1)} \right) \\
&= v_0 \kappa^\alpha \sum_{j=0}^{n+1} \sum_{i=0}^j \binom{j}{i} \frac{(-1)^j (\kappa^\alpha + \eta^\alpha)^{j-i} (\kappa\eta)^{i\alpha} t^{(i+j+1)\alpha}}{\Gamma((i+j+1)\alpha+1)},
\end{aligned}$$

which completes the argument regarding u_n .

Finally, suppose that (24) holds for $n \geq 1$. From (22), (11) and (14), it follows that

$$\begin{aligned}
 v_{n+1}(t) &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + \kappa^{2\alpha} L^{-1} \left[\frac{1}{s^{2\alpha}} L v_n(s) \right] \\
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + \kappa^{2\alpha} L^{-1} \left[\frac{1}{s^{2\alpha}} L \left[v_0 \sum_{i=0}^{2n+1} \frac{(-\kappa^\alpha t^\alpha)^i}{\Gamma(i\alpha + 1)} \right] \right] \\
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + \kappa^{2\alpha} L^{-1} \left[\frac{1}{s^{2\alpha}} v_0 \sum_{i=0}^{2n+1} (-\kappa^\alpha)^i L \left[\frac{t^{i\alpha}}{\Gamma(i\alpha + 1)} \right] \right] \\
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + \kappa^{2\alpha} L^{-1} \left[\frac{1}{s^{2\alpha}} v_0 \sum_{i=0}^{2n+1} (-\kappa^\alpha)^i \frac{1}{s^{i\alpha+1}} \right] \\
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} \right) + v_0 \sum_{i=0}^{2n+1} (-1)^i \kappa^{(i+2)\alpha} L^{-1} \left[\frac{1}{s^{(i+2)\alpha+1}} \right] \\
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} + \sum_{i=0}^{2n+1} (-1)^i \kappa^{(i+2)\alpha} \frac{t^{(i+2)\alpha}}{\Gamma((i+2)\alpha + 1)} \right) \\
 &= v_0 \left(1 - \frac{\kappa^\alpha t^\alpha}{\Gamma(\alpha + 1)} + \sum_{i=0}^{2n+1} \frac{(-\kappa^\alpha t^\alpha)^{(i+2)}}{\Gamma((i+2)\alpha + 1)} \right) \\
 &= v_0 \left(\sum_{i=0}^1 \frac{(-\kappa^\alpha t^\alpha)^i}{\Gamma(i\alpha + 1)} + \sum_{i=2}^{2(n+1)+1} \frac{(-\kappa^\alpha t^\alpha)^i}{\Gamma(i\alpha + 1)} \right) = v_0 \sum_{i=0}^{2(n+1)+1} \frac{(-\kappa^\alpha t^\alpha)^i}{\Gamma(i\alpha + 1)},
 \end{aligned}$$

and the proof is complete. \square

6. Numerical Experiments

With the aim to compare our results with some of those obtained in [1] we will consider

$$\kappa = 0.02873, \quad \eta = 0.08442, \quad v_0 = 4, \quad \text{and} \quad t \in [0, 7].$$

For these values, from [27, Theorem 4.1], we know that u_n converges to the solution u of (19) for any $\alpha \geq 0.88632$ due that the condition

$$\frac{(\kappa^\alpha + \eta^\alpha)7^\alpha}{\Gamma(\alpha + 1)} + \frac{(\kappa\eta)^\alpha 7^{2\alpha}}{\Gamma(2\alpha + 1)} < 1$$

is satisfied. On the other hand, v_n converges to the solution v of (21) for any α since the requirement in [27, Theorem 4.1] is always verified.

Some of the α values within the interval $[0.88632, 1)$ are 0.97 and 0.89; for these values, the approximation errors are tabulated in Tables 1 and 2, respectively. It can be observed that the results obtained by the FVIM are more accurate than those yielded by VIM [1] and VIM [24] in all cases, and that, according to [27, Theorem 4.1], the error improves as the number of iterations increases and as α approaches 1.

Table 1. Approximation errors for the FVIM, VIM [1] and VIM [24] for $\alpha = 0.97$

n	$\alpha = 0.97$					
	$\ u - u_n\ _\infty$	$\ u - \psi_n\ _\infty$	$\ u - \chi_n\ _\infty$	$\ v - v_n\ _\infty$	$\ v - \phi_n\ _\infty$	$\ v - \varphi_n\ _\infty$
2	2.0990e-02	7.1449e-02	9.5971e-02	6.6364e-07	6.6284e-03	1.4004e-02
3	3.8142e-03	1.2195e-02	1.8357e-02	6.3351e-10	3.7802e-04	1.0353e-03
4	5.7492e-04	1.6267e-03	2.4934e-03	3.8058e-13	1.7316e-05	4.7095e-05
5	7.4152e-05	1.7963e-04	2.4163e-04	null	6.6364e-07	7.2062e-07
6	8.3667e-06	1.6987e-05	--	null	2.1882e-08	--
7	8.3964e-07	1.4072e-06	--	null	6.3351e-10	--

Table 2. Approximation errors for the FVIM, VIM [1] and VIM [24] for $\alpha = 0.89$

n	$\alpha = 0.89$					
	$\ u - u_n\ _\infty$	$\ u - v_n\ _\infty$	$\ u - \chi_n\ _\infty$	$\ v - v_n\ _\infty$	$\ v - \phi_n\ _\infty$	$\ v - \varphi_n\ _\infty$
2	4.5438e-02	1.17671e-01	2.75256e-01	3.3643e-06	1.2729e-02	7.3797e-02
3	1.0826e-02	2.5248e-02	8.1245e-02	6.5858e-09	9.8529e-04	1.0465e-02
4	2.1697e-03	4.2964e-03	1.4022e-02	6.5858e-09	6.2456e-05	4.8578e-04
5	3.7654e-04	6.1335e-04	7.6709e-04	7.9936e-15	3.3643e-06	1.0588e-04
6	5.7760e-05	7.5888e-05	--	null	1.5797e-07	--
7	7.9533e-06	8.3130e-06	--	null	6.5858e-09	--

In Tables 3 and 4, we compile the error approximations obtained for $\alpha = 0.72$ and $\alpha = 0.62$, both values considered in [1] but outside the interval $[0.88632, 1)$. In these cases, we observe that the FVIM remains the most efficient method for approximating v , while the VIM [1] performs slightly better than the FVIM for approximating u . The lower precision of the VIM [24] in all these cases is the main reason we have only considered the first five terms provided by this method. Nevertheless, it should be noted that the computational cost of obtaining each iteration with this method is higher than that of the other methods.

Table 3. Approximation errors for the FVIM, VIM [1] and VIM [24] for $\alpha = 0.72$

n	$\alpha = 0.72$					
	$\ u - u_n\ _\infty$	$\ u - v_n\ _\infty$	$\ u - \chi_n\ _\infty$	$\ v - v_n\ _\infty$	$\ v - \phi_n\ _\infty$	$\ v - \varphi_n\ _\infty$
2	2.177e-01	3.159e-01	1.4866	9.0844e-05	4.6935e-02	6.1106e-01
3	9.1415e-02	1.08192e-01	8.55439e-01	7.7956e-07	6.7892e-03	1.9295e-01
4	3.3241e-02	3.0269e-02	2.29284e-01	4.7707e-09	8.3796e-04	1.8086e-02
5	1.0729e-02	7.2970e-03	3.9745e-02	2.2218e-11	9.0844e-05	1.5215e-02
6	3.1276e-03	1.5621e-03	--	8.2601e-14	8.8251e-06	--
7	8.3422e-04	3.0253e-04	--	4.4409e-16	7.7955e-07	--

Table 4. Approximation errors for the FVIM, VIM [1] and VIM [24] for $\alpha = 0.62$

n	$\alpha = 0.62$					
	$\ u - u_n\ _\infty$	$\ u - v_n\ _\infty$	$\ u - \chi_n\ _\infty$	$\ v - v_n\ _\infty$	$\ v - \phi_n\ _\infty$	$\ v - \varphi_n\ _\infty$
2	5.1944e-01	5.3679e-01	3.3947	5.6313e-04	9.5226e-02	1.5580
3	3.0242e-01	2.3864e-01	2.6857	1.1120e-05	1.9538e-02	6.9515e-01
4	1.5517e-01	8.8148e-02	9.3978e-01	1.6453e-07	3.5049e-03	8.6431e-02
5	7.1717e-02	2.8468e-02	2.3613e-01	1.9274e-09	5.6313e-04	1.0027e-01
6	3.0307e-02	8.2782e-03	--	1.8556e-11	8.2406e-05	--
7	1.1851e-02	2.2015e-03	--	1.5143e-13	1.1119e-05	--

7. Conclusions

Some systems of first-order differential equations, such as the one used to model ethanol concentration in blood and stomach (cf. (1)), can be rewritten as a set of higher-order differential equations (cf. (19), (21)). Consequently, any method designed to handle these higher-order equations can be used to obtain numerical approximations for the original system. In this work, we explore this approach for the system (1) and the fractional variational iteration method proposed in [27]. We find that the FVIM is effective for obtaining numerical approximations of (1), particularly demonstrating notable performance in addressing the single-term differential equation involved (see the first line in (1)).

We compare the results obtained using the FVIM with those derived from the variational iteration methods VIM [1] and VIM [24]. In this comparison, the FVIM consistently yields lower approximation errors.

Although the FVIM is originally defined by a recurrence relation, we provide explicit expressions for the approximating terms that depend solely on the iteration order, which simplifies and streamlines its implementation. Furthermore, we enhance the implementation of VIM [1] by providing explicit expressions for the approximating terms that avoid the use of cross-terms.

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